

Closed loop control of QETU for Ground State Preparation

Shivane Dadi, Dhruv Patel, Murphy Yuezhen Niu

Department of Computer Science

Agenda

- Intro: Ground state preparation via QETU
- Simulating noise in QETU
- A closed loop control scheme
- Numerical validation

Intro: Ground state preparation via QETU

- Many quantum algorithms for preparing the ground state of a Hamiltonian (e.g. VQE, QPE, Adiabatic evolution)
 - Each with existing tradeoffs
- **Quantum Eigenvalue Transformation (QETU)** offers a framework with low query complexity and a singular ancilla

$$H_{\text{TFIM}} = - \sum_{j=0}^{n-2} Z_j Z_{j+1} - g \sum_{j=0}^{n-1} X_j$$

Dong, Lin, Tong — PRX Quantum 3, 040305 (2022)

QETU Advantage across the rest of the field

TABLE II. A comparison of the performance of quantum algorithms for ground-state preparation in terms of the query complexity, the query depth, the number of ancilla qubits, and the level of multiqubit control (abbreviated as “MQC” in the table) operations needed. γ is the overlap between the initial guess $|\phi_0\rangle$ and the ground state, Δ is a lower bound of the spectral gap, and $1 - \epsilon$ is the target fidelity. “HE” denotes the Hamiltonian-evolution model (assuming no ancilla qubits) and “BE” the block-encoding model. Here, we assume that an upper bound of the ground-state energy is known (μ in Theorem 6). The algorithm in Ref. [37] (GTC19 in the table) requires precise knowledge of the ground-state energy. We assume that Ref. [1] uses m ancilla qubits and that the high-level MQC operation is due to the block encoding of H .

	Query depth	Query complexity	Number of ancilla qubits	Need MQC?	Input model
This work (Theorem 6)	$\tilde{\mathcal{O}}(\Delta^{-1})$	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-2})$	$\mathcal{O}(1)$	No	HE
This work (Theorem 11)	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-1})$	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-1})$	$\mathcal{O}(1)$	Low	HE
QPE (high confidence) [41–43]	$\tilde{\mathcal{O}}(\Delta^{-1})$	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-2})$	$\mathcal{O}(\text{poly log}(\Delta^{-1}\gamma^{-1}\epsilon^{-1}))$	High	HE
QPE (semiclassical) [44,45]	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-2})$	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-4})$	$\mathcal{O}(1)$	No	HE
GTC19 (Theorem 1) [37]	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-1})$	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-1})$	$\mathcal{O}(\log(\Delta^{-1}) + \log \log(\epsilon^{-1}))$	High	HE
LT20 [1]	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-1})$	$\tilde{\mathcal{O}}(\Delta^{-1}\gamma^{-1})$	m	High	BE

Important to note, there now exist other methods that are more efficient than QETU. but QETU is still a necessary algorithm for downstream tasks

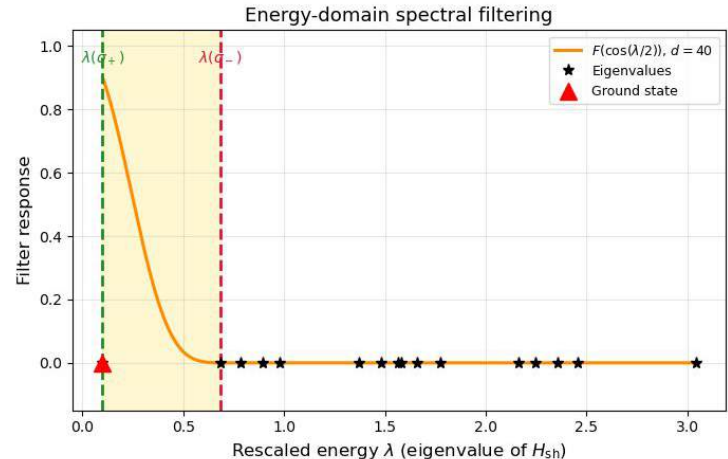
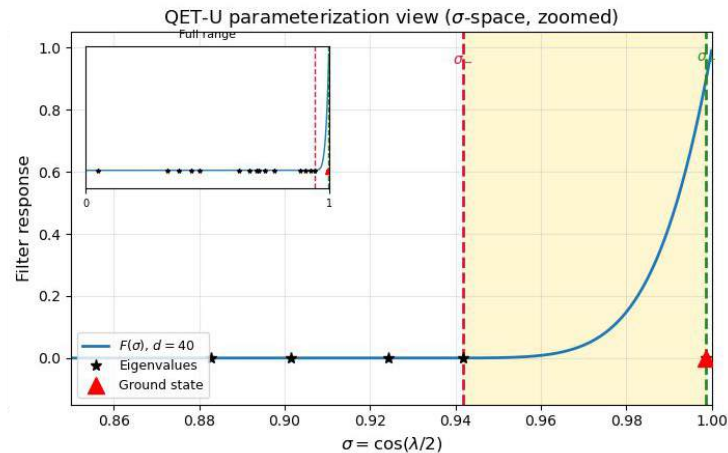
Dong, Lin, Tong — PRX Quantum 3, 040305 (2022)

QETU filters for ground state

- **Goal:** apply a function $F(H)$ to $|\psi_{\text{init}}\rangle$ | such that $F(E_0) \approx 1$ and $F(E_k) \approx 0$ for all excited states

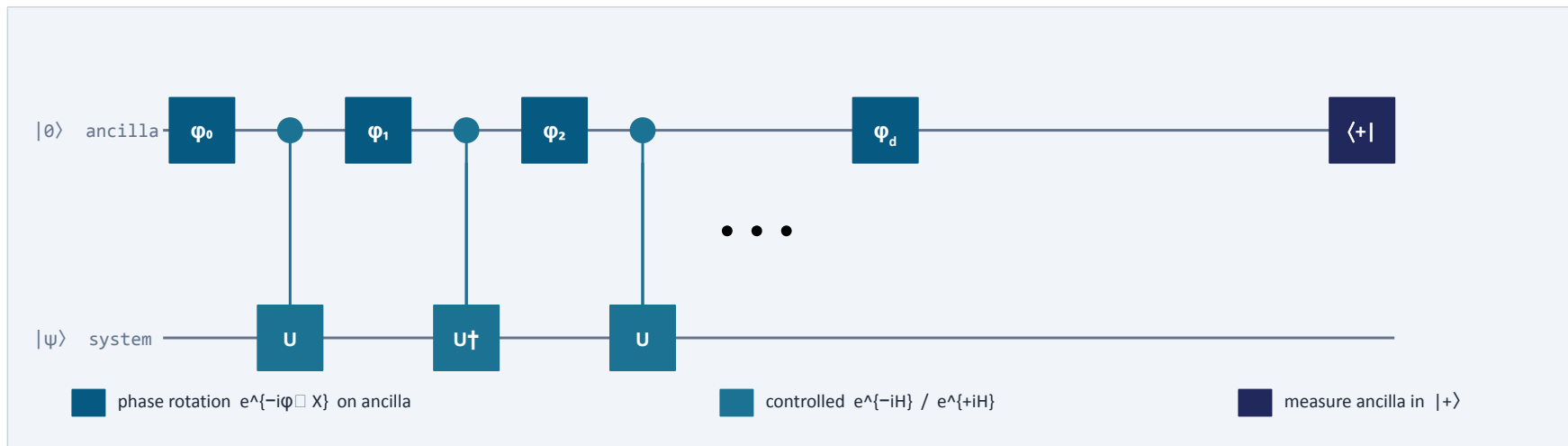
$$F(H) |\psi_{\text{init}}\rangle = c_0 F(E_0) |\psi_0\rangle + \sum_{k \geq 1} c_k F(E_k) |\psi_k\rangle$$

- **QET-U realizes this as a polynomial** $F(\sigma)$ in $\sigma = \cos(\lambda/2)$, where $\lambda \in [0, \pi]$ is the rescaled energy
- Sharper filter $\rightarrow$ better projection



The QETU circuit in detail

The phase angles $\{\phi_j\}$ are computed classically so that the circuit implements the target polynomial $F(\cos(H/2))$



On success (ancilla measured in $|+\rangle$): system register is left in $F(\cos(H/2)) |\psi\rangle$ the designed filter is applied to each eigenstate.

On failure ($|-\rangle$): a different polynomial was applied; discard and rerun.

$$U = e^{-iH_{\text{tfim}}}$$

$$R_x(\phi_j) = e^{-i\phi_j X/2}$$

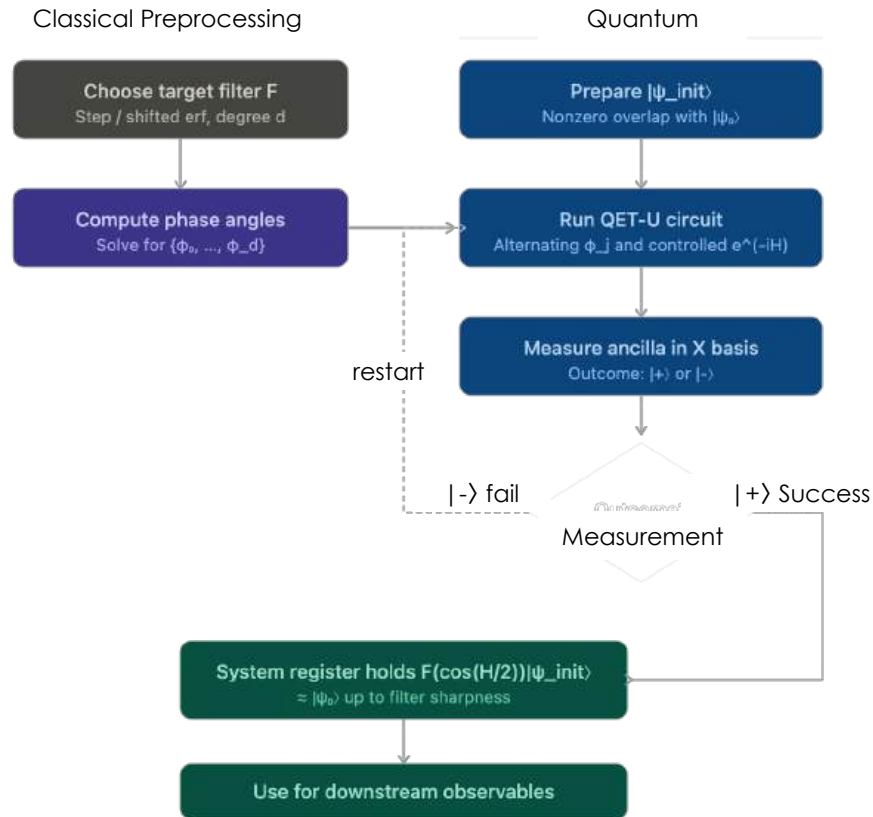
Quantum Signal Processing

- QSP is used to perform polynomial transformation on matrices
- Foundational framework of many quantum algorithms
- Achieves this by applying a sequence of specific rotations, parameterized by **phase factors**

$$U(x, \Phi) = \begin{pmatrix} p(x) & r(x) \\ r^*(x) & p^*(x) \end{pmatrix} = e^{i\phi_0 Z} e^{i \arccos(x) X} e^{i\phi_1 Z} e^{i \arccos(x) X} \dots e^{i\phi_{d-1} Z} e^{i \arccos(x) X} e^{i\phi_d Z},$$

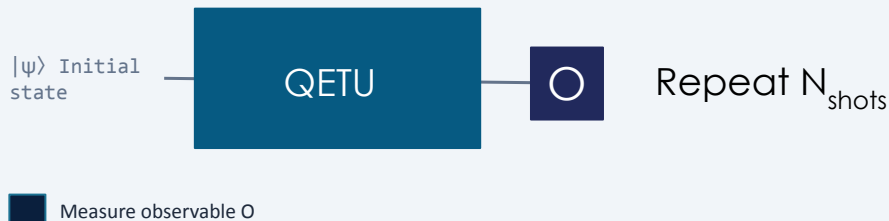
$$U_x \equiv \begin{bmatrix} x & * \\ * & * \end{bmatrix} \Rightarrow U_{f(x)} \equiv \begin{bmatrix} f(x) & * \\ * & * \end{bmatrix}.$$

QETU as a whole



From ground state to observables

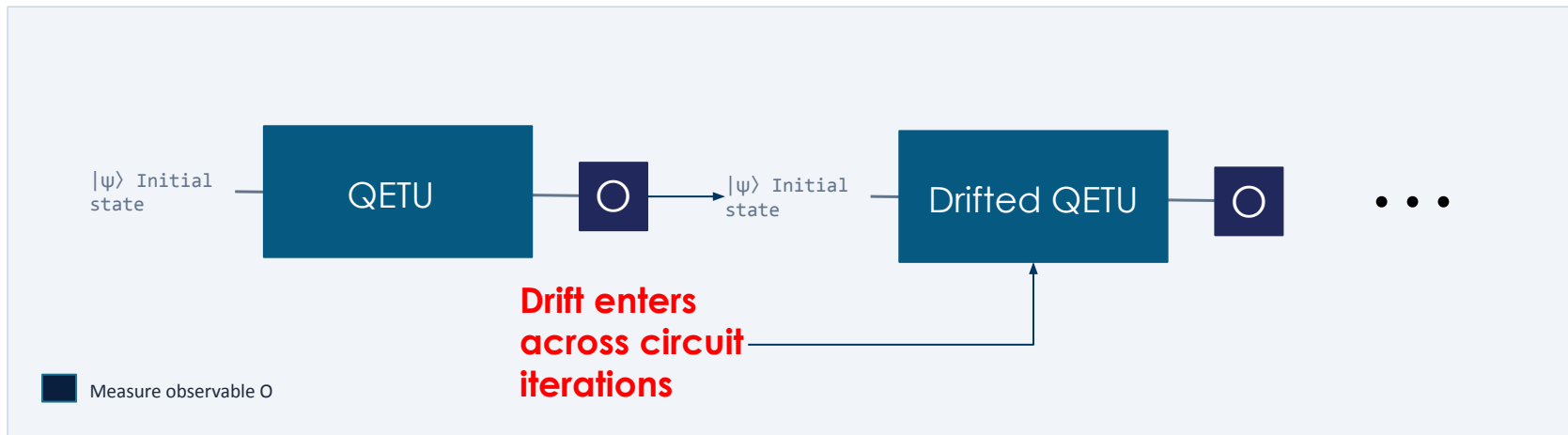
- We now take the prepared ground state and measure on it. Typical observables are $\langle \psi_0 | O | \psi_0 \rangle$
 - Energy ($O = H$), correlation functions, order parameters
- Run QET-U many times $\rightarrow$ collect successes $\rightarrow$ measure O on each $\rightarrow$ average gives $\langle O \rangle \approx \langle \psi_0 | O | \psi_0 \rangle$ up to filter approximation error.
- We study **energy density** $\langle H \rangle / N$, thermodynamic-limit observable



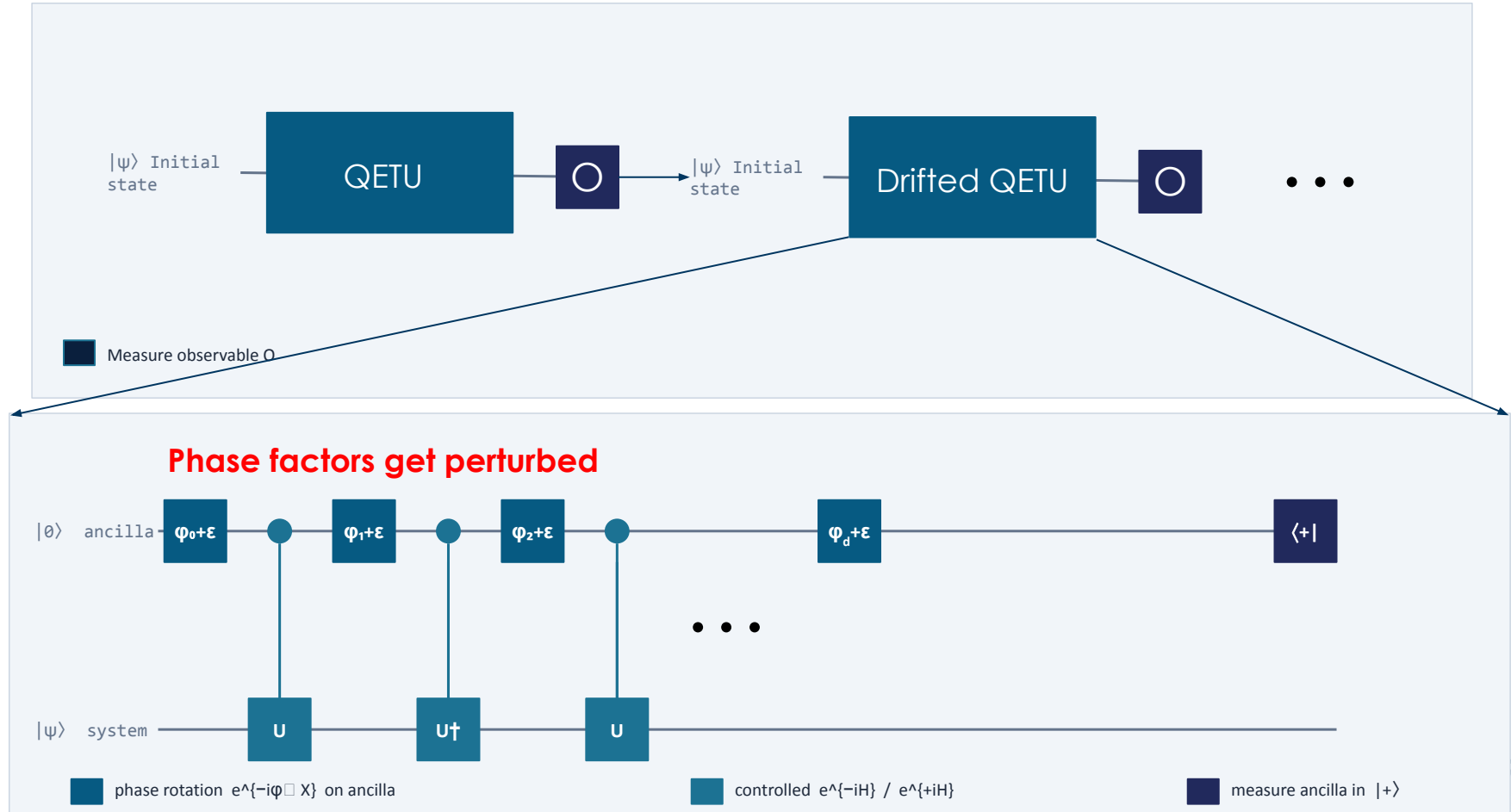
Introducing noise to QETU

Hardware Defects

- QETU makes assumptions about exact implementation of gates
- Real hardware suffers from calibration drift over large timescale
 - Timescales required for measuring observables on the ground state



Hardware Defects



Noise Models

Additive

$$\phi^{(\text{noisy})} = \phi_{W_Z} + \sigma \cdot \frac{j}{d_{\max}}$$

Multiplicative

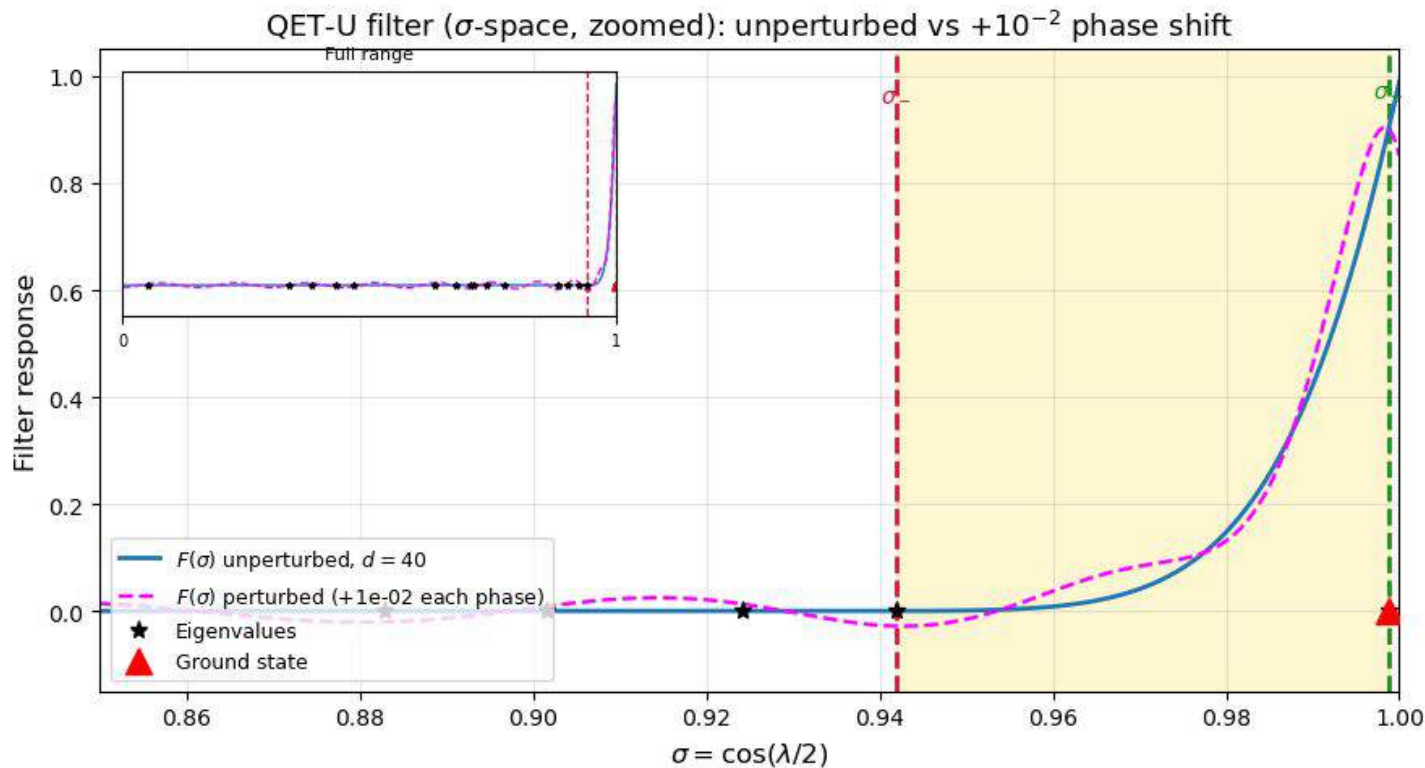
$$\phi^{(\text{noisy})} = \phi_{W_Z} \cdot \left(1 + \sigma \cdot \frac{j}{d_{\max}}\right)$$

σ = Drift Rate

j = Current Circuit Iteration

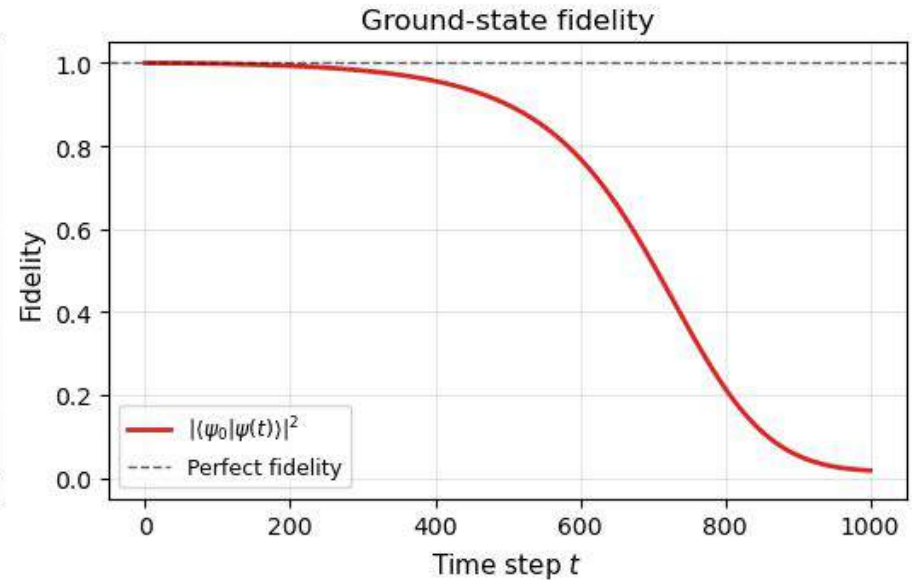
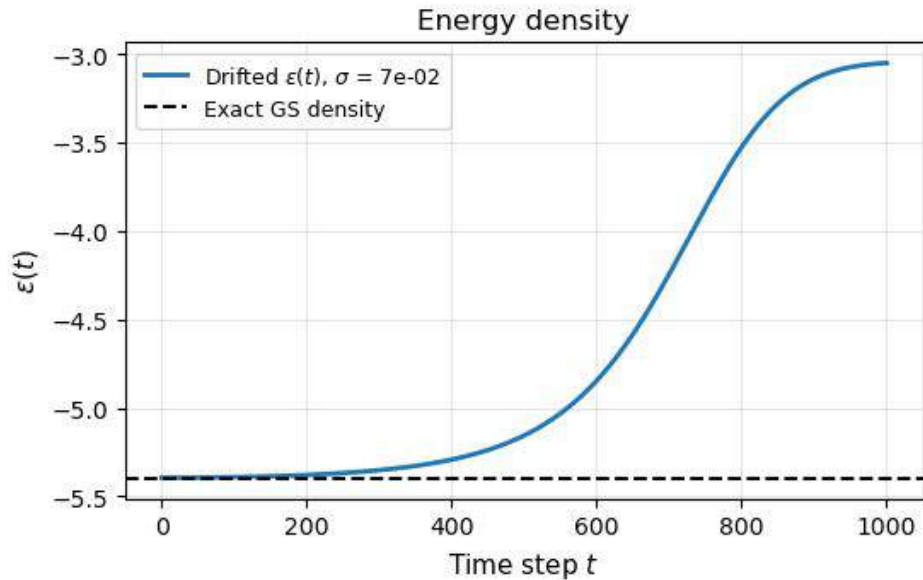
$d_{\max}$ = Max Circuit Iterations

Noise Models



Noise Models

Drifted QET-U preparation (fixed $\sigma = 7\text{e-}02$, $n = 4$, $g = 4.0$)



A Closed Loop Control Scheme

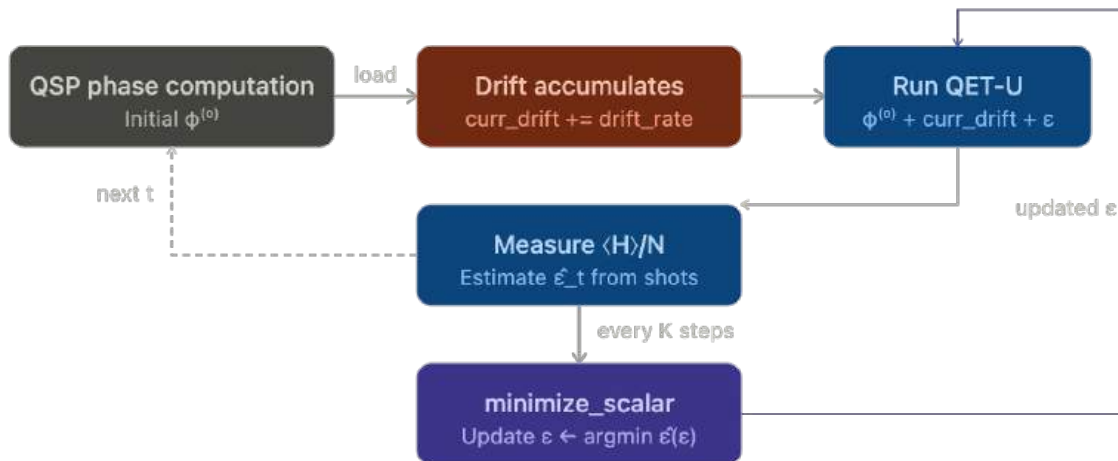
Energy density as a learning signal

- Drift causes the implemented phases $\{\phi_j'\}$ to drift away from the designed $\{\phi_j\}$ → filter degrades → **prepared state has higher energy** than the true ground state.
- This makes $\langle H \rangle / N$ a natural loss function: minimize it by adjusting the phases
- *Closed-loop QET-U: treat phases as learnable parameters, use measured energy density as the objective, optimize online*

Closed loop optimization scheme

Offline (one-time)

Per-step loop (every t until $t_{\max}$)



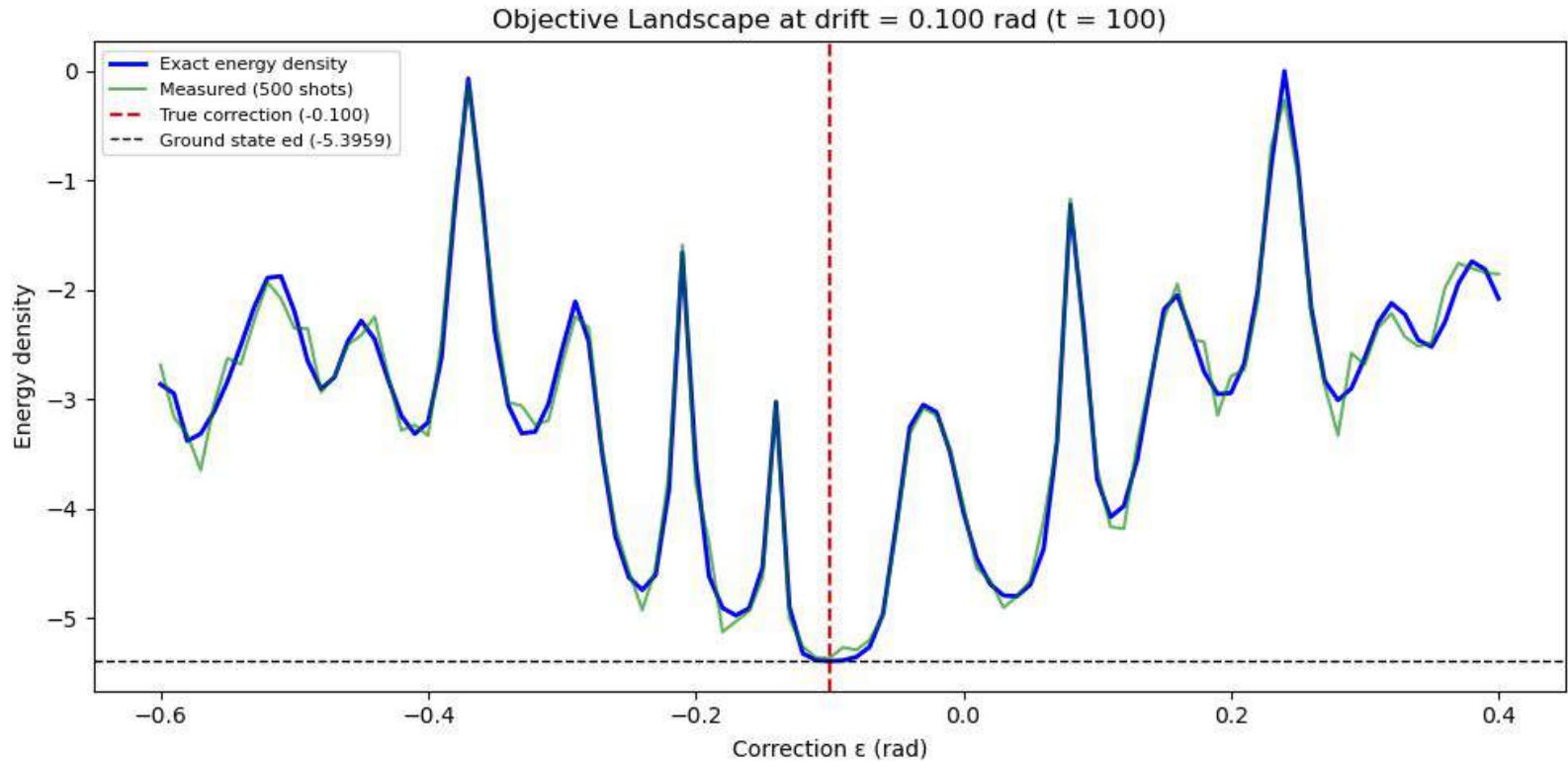
Single-parameter correction

ϵ is one scalar applied uniformly to every phase — not a per- ϕ_j update

Bounded local search

ϵ searched in $[\epsilon_{\text{prev}} - \text{margin}, \epsilon_{\text{prev}} + \text{margin}]$ — avoids oscillatory local minima

Closed loop optimization scheme



Experimental Setup

System & Hamiltonian

TFIM Hamiltonian

$$H = -\sum_j Z_j Z_{j+1} - g \sum_j X_j$$

System size

$n = 4$ qubits

Transverse field

$$g = 4$$

Filter degree

$$d = 40$$

Drift rate

$$\delta_{\text{rate}} \in [10^{-4}, 10^{-2}] \text{ rad/step}$$

Measurement budget

$$N_{\text{total}} = t_{\text{max}} + \left\lfloor \frac{t_{\text{max}}}{K} \right\rfloor \cdot E_{\text{opt}} \cdot M_{\text{shots}}$$

Total shots for experiment

$$N_{\text{total}}$$

Time horizon

t_{max} steps

Recalibration interval

K steps

Optimizer evals per recalibration

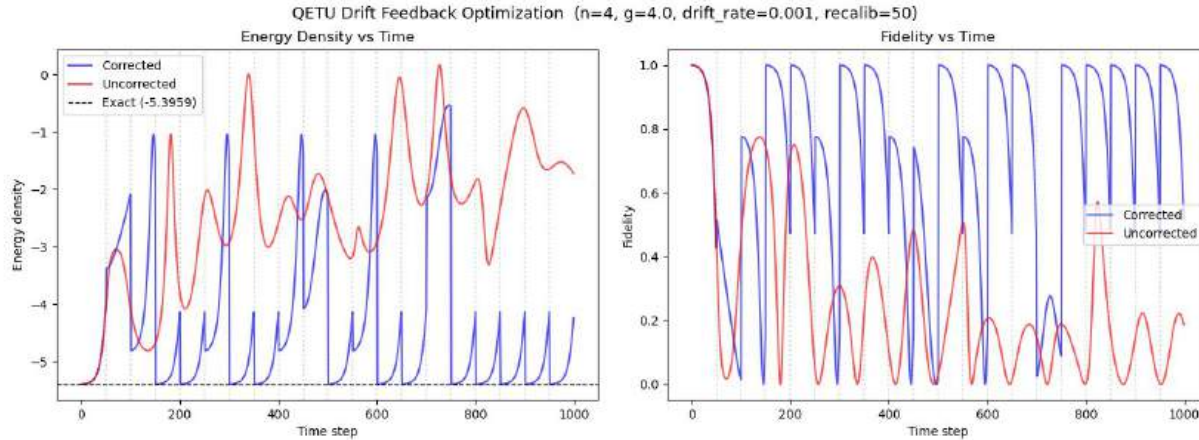
$$E_{\text{opt}}$$

Per-step cost

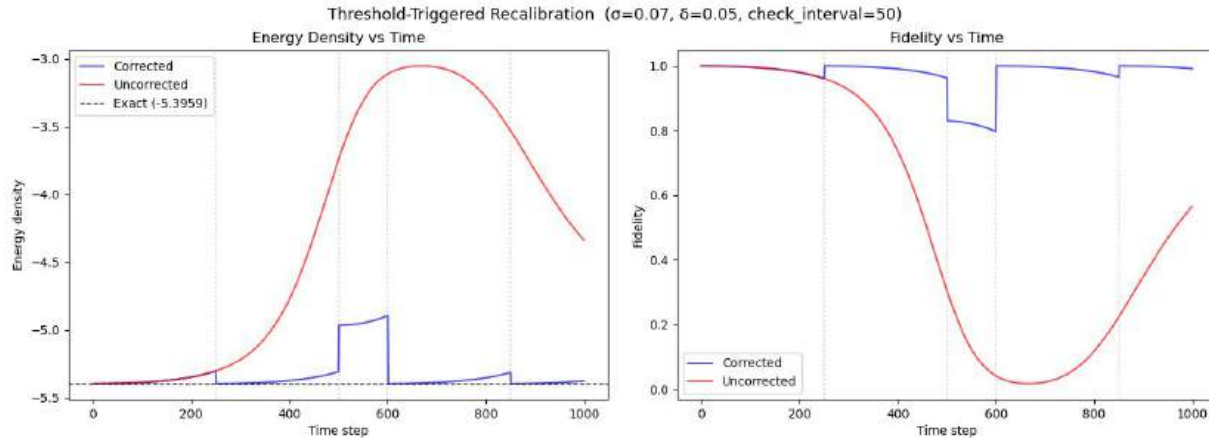
$$M_{\text{shots}}$$

Numerical Validation

Additive

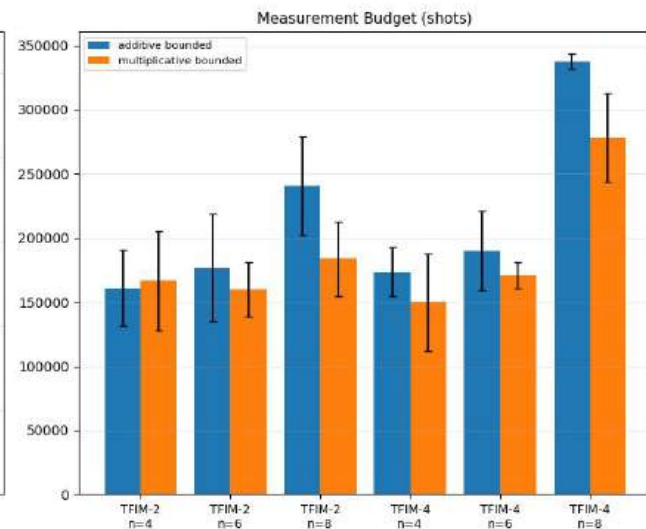
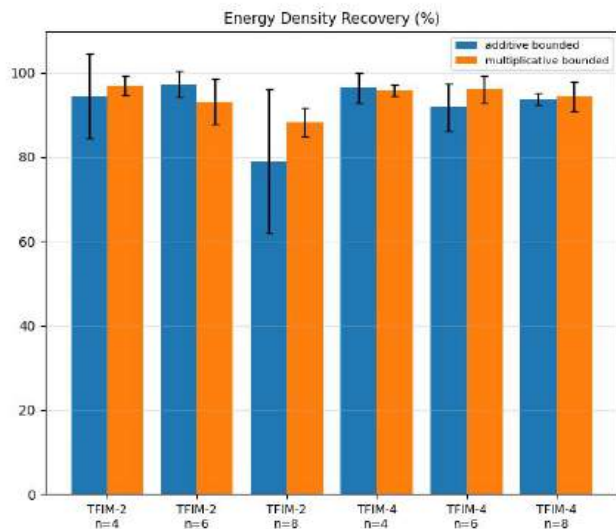
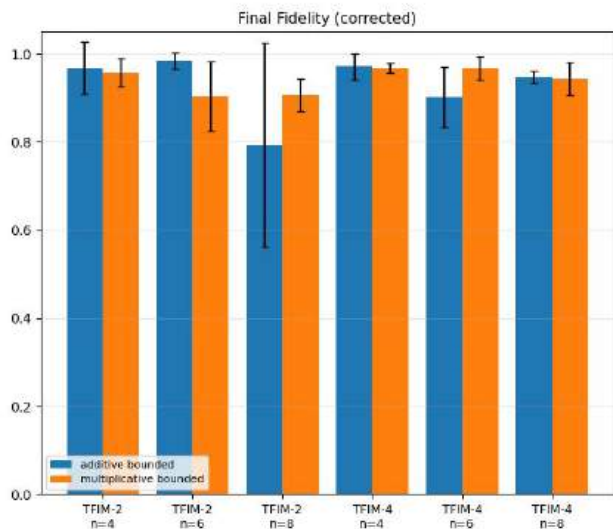


Multiplicative



Numerical Validation - scaling

Drift Correction Performance — Bounded models only (from JSON)



Future Directions

- Continue Scaling to larger system sizes (20+ qubits)
- Test framework under different hamiltonian (XY model)

References

- Dong Y, Lin L, Tong Y. Ground-State Preparation and Energy Estimation on Early Fault-Tolerant Quantum Computers via Quantum Eigenvalue Transformation of Unitary Matrices. PRX quantum. 2022;3(4). doi:<https://doi.org/10.1103/prxquantum.3.040305>
- Ying, L. (2022). Stable factorization for phase factors of quantum signal processing. *Quantum*, 6, 842.
- Kikuchi, Yuta, et al. "Realization of quantum signal processing on a noisy quantum computer." npj Quantum Information 9.1 (2023): 93

UC SANTA BARBARA